

Enlarging the 2D-Brownian filtration with a subordinated perpetuity.

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1)

1. Introduction.

(1) Let $(B_t, t \geq 0)$ and $(Y_t, t \geq 0)$ be two independent 1-dimensional Brownian motions. For $\mu < 0$, and $\gamma \in \mathbb{R}$, we consider the associated Brownian motions with respective drifts μ and γ :

$$B_t^{(\mu)} = B_t + \mu t, \quad \text{and} \quad Y_t^{(\gamma)} = Y_t + \gamma t, \quad t \geq 0,$$

and we define the process: $X_t^{(\mu, \gamma)} = \int_0^t dY_s^{(\gamma)} \exp(B_s^{(\mu)}), \quad t \geq 0.$

It has recently been remarked ([1], [6]) that the law of the so-called subordinated perpetuity:

$$X_{\infty}^{(\mu, \gamma)} \stackrel{\text{def}}{=} X_{\infty}^{(\mu, \gamma)} \quad (\text{which is well defined, since } \mu < 0)$$

is the generic type IV Pearson distribution (Pearson [4], Johnson-Kotz ([2]), Chapter 12), Wong [5], ...); precisely, one has the following

Theorem 1 ([1], [6]): The law of $X_{\infty}^{(\mu, \gamma)}$ admits the

density: (1) $f_{\mu, \gamma}(x) = \frac{C_{\mu, \gamma}}{(1+x^2)^{\frac{1}{2}-\mu}} \exp(2\gamma \arctan(x))$

For simplicity, we shall now write only f for $f_{\mu, \gamma}$; the following quantities will play an important role in the sequel:

If $\gamma = 0$, in particular, $C_{\mu, 0} = ?$
(Via de Haan paper).

$$\left\{ \frac{1}{\int_0^t ds \exp(\alpha B_s^{(\mu)})}, t > 0 \right\} \stackrel{\text{(law)}}{=} \left\{ \frac{1}{\int_0^t ds \exp(\alpha \tilde{B}_s^{(-\mu)})} + \frac{1}{\int_0^\infty ds \exp(\alpha \tilde{B}_s^{(\mu)})}, t > 0 \right\}$$

where, on the RHS,
positive drift $(-\mu)$, and
independent of $B^{(-\mu)}$.
 $\tilde{B}^{(\mu)}$ denotes a Brownian motion with
a copy of $B^{(\mu)}$, assumed to be

(1.2) Our motivation to develop this case study of enlargements come from two origins, at least:

a) there is presently a lot of interest in Mathematical Finance to study how inside trading may modify the pricing framework on a given market ; mathematically, this may be translated in terms of an enlargement of filtration;

b) we hope the present study may help us to develop further extensions of Pitman's theorems just as [3] led us to [3']

2. Proof of Theorem 2.

(to be developed more explicitly)

(2.1) The presentation of the initial enlargement formula given in Chapter 12, p. 33-34 of [7] applies (with only one change, made necessary by the fact that our filtration $\{F_t\}$ is generated by a 2D-Brownian motion instead of a one-dimensional one).

We may summarize this presentation as follows:

denoting $\phi_x(t) = \phi(t, x)$ the density of the conditional distribution of X given F_t , and writing:

$$\phi_x(t) = \phi(0) \exp \left\{ \int_0^t (\rho_1(s, x) dB_s + \rho_2(s, x) dY_s) - \frac{1}{2} \int_0^t ds [\rho_1^2 + \rho_2^2](s, x) \right\}$$

a generic $\{F_t\}$ martingale :

$$M_t = \int_0^t (m_1(s) dB_s + m_2(s) dY_s)$$

$$\left\{ \frac{1}{\int_0^t ds \exp(\alpha B_s^{(\mu)})}, t > 0 \right\} \stackrel{\text{(law)}}{=} \left\{ \frac{1}{\int_0^t ds \exp(\alpha \tilde{B}_s^{(-\mu)})} + \frac{1}{\int_0^\infty ds \exp(\alpha \tilde{B}_s^{(\mu)})}, t > 0 \right\}$$

where, on the RHS, $\tilde{B}^{(-\mu)}$ denotes a Brownian motion with positive drift $(-\mu)$, and $\tilde{B}^{(\mu)}$ independent of $\tilde{B}^{(-\mu)}$.

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We may summarize this presentation as follows:

denoting $\phi_x(t) = \phi(t, x)$ the density of the conditional distribution of X given \mathcal{F}_t , and writing:

$$\phi_x(t) = \phi(0) \exp \left\{ \int_0^t (\rho_1(s, x) dB_s + \rho_2(s, x) dY_s) - \frac{1}{2} \int_0^t ds [\rho_1^2 + \rho_2^2](s, x) \right\}$$

a generic $\{\mathcal{F}_t\}$ martingale :

$$M_t = \int_0^t (m_1(s) dB_s + m_2(s) dY_s)$$

is decomposed in $\{\hat{F}_t\}$ as:

$$M_t = \hat{M}_t + \int_0^t ds \left\{ m_1(s) \rho_1(s, X) + m_2(s) \rho_2(s, X) \right\}$$

thus, in order to prove theorem 2, it only remains to find ρ_1 and ρ_2 .

(2.2) the process $\{\phi_x(t); t \geq 0, x\}$ of conditional densities given $\{\hat{F}_t\}$ is now found to be (with the help of theorem 1):

$$\boxed{\phi_x(t) = \frac{1}{e^{(p)}_t} f\left(\frac{x - X_t}{e^{(p)}_t}\right)}, \text{ where: } e^{(p)}_t = \exp(B^{(p)}_t).$$

Note, in particular, that this provides an interesting family of $\{\hat{F}_t\}$ martingales.

Next, ρ_1 and ρ_2 are obtained after writing $\{\phi_x(t), t \geq 0\}$ in exponential form; one finds:

$$\rho_1(s, x) = -\varphi\left(\frac{x - X_s}{e^{(p)}_s}\right), \text{ and } \rho_2(s, x) = -\varphi\left(\frac{x - X_s}{e^{(p)}_s}\right),$$

which completes the proof of theorem 2.

3. A discussion of the enlargement formula (3).

To start with, we would like to compare formula (3) with the (enlargement) formula for the filtration of $(B_t, t \geq 0)$, enlarged with:

$$A^{(p)} = A^{(p)}_\infty = \int_0^\infty ds \exp(2B^{(p)}_s),$$

which is presented in [3], where it reads as follows:

$$(5) \quad B^{(p)}_t = B^{*(-p)}_t - \int_0^t ds \frac{\exp(2B^{(p)}_s)}{(A^{(p)} - A^{(p)}_s)},$$

with the notation $\{B^{*(-p)}_t\}$ for an $\mathcal{F}_t^* = \mathcal{F}_t \vee \sigma\{A^{(p)}\}$ Brownian motion.

(at this point, we warn the reader that in [3], the negative drift is denoted by $(-\mu)$, so that μ in [3] is changed in $(-\mu)$ here). We now make several comparisons between formulae (3) and (5).

(3.1) Let us consider the case $\gamma=0$ in Theorem 2, so that the function ψ becomes:

$$\psi(z) = (2\mu) + \frac{1-2\mu}{1+z^2}$$

Let us further remark that formula (5) is also the enlargement formula in $\mathcal{F}_t^* = \mathcal{F}_t \vee \sigma(A^{(\mu)}) \vee \sigma(\hat{\gamma}_u, u \geq 0)$, where $(\hat{\gamma}_u, u \geq 0)$ denotes the Dubins-Schwarz Brownian motion associated with the martingale:

$$X_t^{(\mu)} = \int_0^t d\hat{\gamma}_s \exp(B_s^{(\mu)}) = \hat{\gamma}_{A_t^{(\mu)}}.$$

Since $\hat{\mathcal{F}}_t = \mathcal{F}_t \vee \sigma(X_\infty^{(\mu)}) \subset \mathcal{F}_t^*$, for every t , in order that formulae (5) and (3) be coherent, the following conditional expectation relation must hold:

$$E \left[\frac{\exp(2B_s^{(\mu)})}{A_s^{(\mu)} - A_s^{(\mu)}} \mid \hat{\mathcal{F}}_A \right] = \frac{(1-2\mu)}{1 + \left(\frac{X-X_A}{e_A^{(\mu)}} \right)^2}.$$

or equivalently *, since $(e_A^{(\mu)})^2 = \exp(2B_A^{(\mu)})$ is $\hat{\mathcal{F}}_A$ measurable

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$$E \left[\frac{1}{A_s^{(\mu)} - A_s^{(\mu)}} \mid \hat{\mathcal{F}}_A \right] = \frac{(1-2\mu)}{(e_A^{(\mu)})^2 + (X-X_A)^2}.$$

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This relationship should follow from the known facts:

conditionally on \mathcal{F}_0 , the pair $(X-X_A, A-A_A)$ is distributed as

$(e_A^{(\mu)} \sqrt{H} N, (e_A^{(\mu)})^2 H)$, where $H \stackrel{\text{(law)}}{=} \frac{1}{2\mu}$, and N is $\mathcal{N}(0,1)$

independent from H .

donc la loi nous dit H

$$\boxed{\text{Partie } 2 \text{ de } A^{(\mu)} \text{ et } A^{(\mu),1}}$$

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