

7 Juillet.

Quelques formules de grossissement initial (1^u - expart).

La formule générale (Thm 12.1, p. 33-34, Z II):

Si $\tilde{\lambda}_t(d\ell) = \lambda_t(d\ell) \rho(\ell, t)$, alors si (X_t) est une (\mathcal{F}_t) martingale, on a: $X_t = \tilde{X}_t + \int_0^t \rho(L, s) d\langle X, B \rangle_s$.

Exemple 1: $L = \int_0^\infty \varphi(s) dB_s$; $\int_0^\infty ds \varphi^2(s) < \infty$.

Alors, λ_t , et $\tilde{\lambda}_t$ sont faciles à calculer, et on trouve:

$$\rho(\ell, s) = \varphi(s) \frac{\ell - m_s}{\sigma_s^2}, \quad \text{où: } m_s = \int_0^s \varphi(u) dB_u; \\ \sigma_s^2 = \int_0^s \varphi^2(u) du.$$

Cas particulier: $\varphi(s) = 1_{[0, t_0]}(s)$; i.e.: $L = B_{t_0}$.

On obtient ainsi: $\rho(\ell, s) = \frac{\ell - B_s}{(t_0 - s)}$, $s < t_0$.

et on a donc: $B_t = \tilde{B}_t + \int_0^{t \wedge t_0} ds \frac{B_{t_0} - B_s}{t_0 - s}$.

Comme maintenant (\tilde{B}_t) est indépendant de $\sigma(B_{t_0})$, on peut conditionner \mathcal{F}_t sur $B_{t_0} = y$ (par ex: 0), et on trouve ainsi la décomp. du pont:

$$B_t = y + \tilde{B}_t + \int_0^t ds \frac{y - B_s}{(t_0 - s)}$$

À la suite, on se servira bcp. de la conséquence suivante pour le pont de Bessel :

$$r(t) = r + \rho t + \frac{n-1}{2} \int_0^t \frac{ds}{r(s)} - \int_0^t \frac{ds}{(t_0-s)} r(s).$$

En fait, ceci nous mène au second Exemple :

Exemple 2 : grossissement de la filtration d'un processus de Markov avec X_{t_0} .

Sur un espace canonique, ou avec \mathcal{F}_u (\mathcal{F}_u) mesurable, on a :

$$(1) \quad E_x [F_u | X_1 = y] = E_x \left[F_u \frac{p_{1-u}(X_u, y)}{p_1(x, y)} \right]$$

et donc, une (\mathcal{F}_u) martingale (M_u) devient, pour la loi du pont : $\mathbb{P}_{x \rightarrow y}^1$,

$$M_u = \tilde{M}_u + \int_0^u \frac{d \langle p_{1-u}(X_u, y), M \rangle_u}{p_{1-u}(X_u, y)}$$

Par ex, pour un mart. b. avec drift : $X_t = B_t + \int_0^t du b(X_u)$, on obtient :

$$B_t = \tilde{B}_t + \int_0^t \frac{\partial}{\partial x} (\log p_{1-u}(X_u, y)) du,$$

A donc, pour le pont : $\mathbb{P}_{x \rightarrow y}^1$, le drift b est changé en :

$$b_y(x) \equiv b(x) + \frac{\partial}{\partial x} (\log p_{1-u}(x, y)).$$

Démonstration de la formule (1) :

$$E_x [f(X_1) F_u] = E_x [f(X_1) E_x [F_u | X_1]]$$

$$= E_x [P_{1-u} f(X_u) F_u] = \int dy p_1(x, y) f(y) E_x [F_u | X_1 = y]$$

$$= E_x [F_u \int dy p_{1-u}(X_u, y) f(y)]$$

Exemple 3:

Décomposition du mouvement brownien

$\{B_u, u \leq d_1\}$ sur $[0, q_1], [q_1, d_1], [q_1, 1], \dots$

- On commence par introduire l'op. de Brownian scaling;
- Puis, la décomposition: (faire un dessin).

$$b_u \equiv \frac{1}{\sqrt{q_1}} B_{uq_1}, u \leq 1, \quad m_u \equiv \frac{1}{\sqrt{t-q_1}} |B_{q_1+u(t-q_1)}|, u \leq 1;$$

$$e_u \equiv r_u \equiv \frac{1}{\sqrt{t-q_1}} |B_{q_1+u(t-q_1)}|, u \leq 1.$$

Cet exemple 3 sera décomposé en 2 parties:

(3.1) grossissement initial avec $L = q = q_1$.

Les formules pour $\lambda_s(dt)$ et $\dot{\lambda}_s(dt)$ sont: $\equiv \equiv \equiv e(s, l)$

$$\lambda_s(dt) \equiv \varepsilon_{q_1}(dt) \Phi\left(\frac{|B_s|}{\sqrt{1-s}}\right) + \frac{dt}{\pi} \int_{s < l < 1} \exp\left(-\frac{B_s^2}{2(l-s)}\right)$$

d'où l'on déduit (mais cela utilise la formule de balayage; à démontrer)

$$\dot{\lambda}_s(dt) \equiv \varepsilon_{q_1}(dt) \Phi'\left(\frac{|B_s|}{\sqrt{1-s}}\right) \frac{\text{sgn}(B_s)}{\sqrt{1-s}} + dt e(s, l) \left(-\frac{B_s}{(l-s)}\right)$$

En conséquence, on a donc:

$$p(t, s) = \left(\frac{\Phi'}{\Phi}\right)\left(\frac{|B_s|}{\sqrt{1-s}}\right) \frac{\text{sgn}(B_s)}{\sqrt{1-s}} \mathbb{1}_{\left(\frac{q_1}{t} \leq s\right)} + \mathbb{1}_{(t > s)} \left(-\frac{B_s}{l-s}\right) \quad (5)$$

La formule de Girsanov devient donc :

$$B_t = \tilde{B}_t + \int_0^t ds \left\{ 1_{(s < q)} \left(-\frac{B_s}{(q-s)} + 1_{(q \leq s)} \left(\frac{\Phi'}{\Phi} \right) \left(\frac{|B_s|}{\sqrt{1-s}} \right) \right) \right.$$

En d'autres termes, on trouve bien le pont brownien $(B_u, u \leq q)$, et le meandre pour $t \geq q$.

A ce point, il faut présenter la relation d'Imhof pour le meandre, i.e.

$$(3) \quad M = \frac{c}{X_1} \cdot P_0^{(3)}$$

En effet, ce qui se passe est que : $E_0^{(3)} \left(\frac{1}{X_1} \middle| \mathcal{F}_t \right) = \psi \left(\frac{X_t}{\sqrt{1-t}} \right)$,

$$\text{or : } \psi(x) = \frac{\Phi(x)}{x}$$

Alors, lorsque l'on écrit la formule de Girsanov pour le changt. en (3), on trouve bien :

$$(4) \quad X_t = B_t + \int_0^t \frac{ds}{X_s} + \int_0^t \frac{\psi'}{\psi} \left(\frac{X_s}{\sqrt{1-s}} \right) \frac{ds}{\sqrt{1-s}}$$

Or, $\frac{\psi'}{\psi}(x) = \frac{\Phi'}{\Phi} - \frac{1}{x}$, et donc de (4), on déduit bien que la formule de Girsanov ci dessus est en accord.....

(3.2) Girsanov initial avec $L = d \equiv d_1$.

Dans ce cas, on trouve :

$$(5) \quad \lambda_s(dl) = \left(\frac{|x|}{\sqrt{2\pi}(l-s)^3} \exp\left(-\frac{x^2}{2(l-s)}\right) \right) \Big|_{x=B_s} 1_{(l \geq s)} dl.$$

On remarque bien sûr que ceci est :

$$\frac{y \exp(-y^2/2v)}{\sqrt{2\pi v^3}} \equiv -\frac{\partial}{\partial y} \left(\frac{1}{v} p(y) \right).$$

On a donc obtenu : $p(t, s) = \theta(v, x) \Big|_{\substack{x=B_t \\ v=t-s}}$ $1(l \geq s)$,

$$\text{avec : } \theta(v, x) = \frac{\frac{\partial^2}{\partial y^2} p(y)}{\frac{\partial}{\partial y} p(y)} \equiv \frac{1}{x} - \frac{x}{v} \quad (x=y!)$$

Finalement, la formule de grossissement donne :

$$(6) \quad B_t = B_1 + \int_1^{t \wedge d} du \left\{ \frac{1}{B_u} - \frac{\frac{1}{2} B_u}{(d-u)} \right\},$$

ce qui montre que le processus $\left\{ B_{1+v}; v \leq d-1 \right\}$

est un pont de Bessel 3, de longueur $(d-1)$.

Démonstration des formules (2) et (5) :

Exemple 4: élargissement de la filtration brownienne avec:

$$L \equiv A_{\infty}^{(\nu)} = \int_0^{\infty} ds \exp 2(B_s + \nu s), \quad \nu < 0; [\mu$$

Pour le calcul de $\lambda_s^{(\nu)}(dt)$, il suffit seulement à l'aide de la connaissance de la loi de $A_{\infty}^{(\nu)}$, lorsque $B_0 = 0$, i.e.:

$$P(A_{\infty}^{(\nu)} \in dy) = P\left(\frac{1}{2Z_{\mu}} \in dy\right).$$

d'écrire:

$$\begin{aligned} \int \lambda_s^{(\nu)}(dt) f(t) &= E[f(A_{\infty}^{(\nu)}) | \mathcal{F}_s] = E[f(A_s^{(\nu)} + A_{\infty}^{(\nu)} \circ \theta_s) | \mathcal{F}_s] \\ &= E_{B_s^{(\nu)}} [f(A_s^{(\nu)}(\omega) + \tilde{A}_{\infty}^{(\nu)})]. \end{aligned}$$

On trouve alors, assez facilement, en posant:

$$A_t^{(\nu)} \equiv \int_0^t ds \exp 2(B_s + \nu s); \quad r_t = \exp(B_t + \nu t):$$

$$p(s, t) \equiv r_t \frac{\partial}{\partial r} (\log \Phi)(r_t, A_t^{(\nu)}; t), \quad [\mu =$$

où :

$$\Phi(r, a; l) = \left(\frac{r^2}{2}\right)^\mu \frac{1}{(l-a)^{\mu+1}} \exp\left(-\frac{r^2}{2(l-a)}\right)$$

En conséquence,

$$\rho(s, l) = (2\mu) - \frac{r^2}{l - A_s^{(\nu)}} ,$$

la formule de grossissement pour le mt. brownien devient :

$$B_t \equiv \beta_t + \int_0^t ds \left\{ (2\mu) - \frac{\exp 2(B_s + \gamma s)}{(A_\infty^{(\nu)} - A_s^{(\nu)})} \right\}$$

$$\equiv \beta_t + \int_0^t ds \left\{ -(2\gamma) - \left(\exp 2(B_s + \gamma s) \right) / \left(\int_1^\infty du \exp 2(B_u + \gamma u) \right) \right\}$$

Vérifications:

a) On peut s'assurer que l'espérance conditionnelle donne bien 0.

$$\text{i.e.} \quad -(2\gamma) - E\left(\frac{1}{A_\infty^{(\nu)}}\right) = -2\gamma - E[2Z_{(-\nu)}] = -2\gamma + 2\gamma = 0.$$

b) Transformons le résultat précédent en un résultat sur les processus de Bond, au moyen de la rep. de Lamperti, i.e.:

$$\exp(B_t + \gamma t) \equiv R \int_0^t ds \exp 2(B_s + \gamma s)$$

Le résultat précédent revient à grossir la filtration de $R^{(\nu)}$ avec $T_0^{(\nu)} = \inf\{t: R_t^{(\nu)} = 0\}$
 Ceci nous donne maintenant :

$$R_u^{(\nu)} \equiv 1 + \gamma u + \frac{2\mu+1}{2} \int_0^u \frac{ds}{R_s^{(\nu)}} - \int_0^u ds \frac{R_s^{(\nu)}}{(T_0 - s)}$$

Autrement dit, $(R_u^{(\nu)}, u \leq T_0)$ est un pont de Bessel d'indice ν issu de 1, allant en 0, par l'intervalle T_0 .

Exemple 5: Yonnest. de la ~~let~~ fct. brownienne avec la tribu du temps local.

Thm 12.2:
$$B_t = B_t^{\text{loc}} + \int_0^t ds \left\{ \frac{1}{B_s} - \frac{B_s}{(B_s - 1)} \right\}. \quad (\text{p. 38-39-40}).$$

$\left. \begin{array}{l} \nu \\ t=0 \end{array} \right\}.$

8th of July.

Second Lecture: Progressive enlargement / Martingales vanishing on the zero set

II. Martingales vanishing on the zero set. $Z = \{(t, \omega) ; t \leq 1, B_t(\omega) = 0\}$

Main thm: To $X \in L^1(\mathcal{F}_1)$, we associate: $X_t = E(X | \mathcal{F}_t)$.

Then, the following are equivalent:

- 1) $X \in \mathcal{M}_1^0$; 2) $E(X | \mathcal{F}_q) = 0$; 3) $X_q = 0$.

II.1. Preparations.

a) The balayage formula: It will give us a nb. of examples of elements of \mathcal{M}_1^0 , but its applicability is much wider.

Proposition: If (X_t) is a semimartingale which vanishes on Z , then:
$$z_{q,t} X_t = \int_0^t z_{q,s} dX_s.$$

In particular, if $(X_t) \in \mathcal{M}_1^0$, then: $(z_{q,t} X_t)$ is also in \mathcal{M}_1^0 .

Other examples: $f(t) |B_t| = \int_0^t f(t_s) d|B_s|$

Hence, $f(t) |B_t| - F(t)$ is a martingale.

Lemma 1: $E[z_{q,1}] = E\left[\int_0^1 dL_u z_u\right]$, where: $dL_u = \frac{du}{\sqrt{1-u}} \sqrt{\frac{2}{\pi}}$

Proof: We have already seen this result, see above, and also in the 1st lecture. But, it can also be understood as a consequence of the balayage formula:

$$E[z_{q,1} | B_1 |] = E\left[\int_0^1 dL_u z_u\right],$$

and: $E[|B_1| | \mathcal{F}_q] = \sqrt{1-q_1} E[m_1] / \text{etc.}$

zero set
of BM.

We shall see more such results later... (3rd lecture, e.g.)

Lemma 2:
$$E[Z_g X_g] = E\left[X \int_0^1 dL_u Z_u\right] \stackrel{!}{=} E\left[E(X|\mathcal{F}_g) \int_0^1 dL_u\right]$$

Proof: the 1st equality follows from lemma 1, and integration by parts.

Question: Que se passe-t-il si (X_t) est une semimartingale? \int par exemple:

$$E[Z_{g,t} X_{g,t}] = \text{etc... ?}$$

II.2. Proof of the theorem. (p. 67).

II.3. Some examples of elements of \mathcal{M}^0 ; a general representation theorem.

We now know that elements of \mathcal{M}^0 are martingales whose terminal value ($\in \mathcal{F}_g$ say) is in the orthogonal of $L^2(\mathcal{F}_g)$, i.e.:

$$Y = X - E[X|\mathcal{F}_g].$$

This gives a systematic way to construct elements of \mathcal{M}_1^0 ;
example:

$$\begin{aligned} X &= f(B_1); & E[X|\mathcal{F}_g] &= E\left[f(\sqrt{1-g} \cdot m_1 \cdot \varepsilon) \mid \mathcal{F}_g\right] \\ &\stackrel{!}{=} & &= \frac{1}{2} \int_0^{\infty} dy \frac{y}{(1-g)} \exp\left(-\frac{y^2}{2(1-g)}\right) (f(y) - f(-y)) \end{aligned}$$

$\frac{1}{2}$ puis, on fait le calcul de:

$$Y_t = P_{1,t} f(B_t) - \frac{1}{2} \int_0^{\infty} dy y (f(y) + f(-y)) E[-$$

or, on a déjà calculé ce genre de quantité.....

Puis, la discussion amène aux questions d'équation conditionnelle.....

III. Some remarks about \mathcal{F}_{L-} and \mathcal{F}_{L+} .

By definition: $\mathcal{F}_{L-} = \sigma\{Z_L; Z(\mathcal{F}_t) \text{ predictable}\}$

$\mathcal{F}_{L+} = \sigma\{Z_L; Z(\mathcal{F}_t) \text{ prog. measurable}\}$

(\mathcal{F}_t) prog. measurable sets/processes..... are complicated, but.....]

Now, since L is a (\mathcal{F}_t) stopping time, we have also the more classical:
etc... \mathcal{F}_{L-} , and:

$$\mathcal{F}_{L+} = (\mathcal{F}^{\text{prog}, L})_L; \quad \mathcal{F}_{L-} = \mathcal{F}_{L-}^{\text{prog}, L}$$

How can we compute cond. expectations w.r. to \mathcal{F}_{L-} and/or \mathcal{F}_{L+} ?

$$E[X | \mathcal{F}_{L-}] = \lim_{u \uparrow L} \frac{E[X 1_{(u < L)} | \mathcal{F}_u]}{P(L > u | \mathcal{F}_u)}$$

$$E[X | \mathcal{F}_{L+}] = \lim_{v \downarrow 0} \frac{E[X 1_{(L < v)} | \mathcal{F}_v]}{(1 - Z_{L+v}^L)}$$

Is \mathcal{F}_{L+} diff. from \mathcal{F}_{L-} ?

Theorem: Let $L = \sup\{t: M_t = 0\}$, where (M_t) is unif. integrable, and $P(M_\infty = 0) = 0$. Then,

\mathcal{F}_{L+} strictly contains \mathcal{F}_{L-} ; in fact, $\text{sgn}(M_\infty)$ is \mathcal{F}_{L+} measurable, and not in \mathcal{F}_{L-} .

Proof: $E[M_\infty | \mathcal{F}_{L-}] = 0$; but, $E[|M_\infty| | \mathcal{F}_{L-}] \neq 0$

In the particular case of Brownian motion, we can show: $\mathcal{F}_{L+} = \mathcal{F}_L \vee \sigma(\text{sgn}(B_L))$

I. Formules de grossissement progressif.

Defn. d'un temps honnête: $L \equiv L_t$, sur $(L \leq t)$.

Fins d'ensembles prévisibles: $L \equiv \sup \{ u : (u, \omega) \in T \}$.

Si L est honnête, alors on définit $\mathcal{F}_t^{\text{prog } L} \equiv \mathcal{F}_t^L$ pour l'instant / Bien sûr attention, mais je préciserai.....

Lemma 1: Si (H_u) est un proc. prévisible \mathcal{F}_t^L , alors, $H_u \equiv H_u' 1_{(u \leq L)} + H_u'' 1_{(u > L)}$
avec deux processus (\mathcal{F}_t) prévisibles H' et H'' .

⊗ Ceci va nous permettre d'obtenir la formule de décomposition, i.e.:
à L , on associe: $Z_t^L = P(L > t | \mathcal{F}_t) \equiv M_t^L - A_t^L$.

Théorème: $X_t = \tilde{X}_t + \int_0^{t \wedge L} \frac{d\langle X, Z \rangle_u}{Z_u^L} + \int_L^t \frac{d\langle X, 1 - Z \rangle_u}{(1 - Z_u^L)}$

To prove the theorem, we first need an easy: / rappel: $Z_t^L = M_t^L - A_t^L$

Lemma 2: a) For every ≥ 0 , prévisible process (Z_u) , one has:

$$E[Z_L] = E\left[\int_0^\infty dA_u^L Z_u\right]$$

b) For every uniformly integrable martingale (X_t) , one has

$$E[X_L] = E[X_\infty A_\infty^L] = E[X_\infty M_\infty^L].$$

$\mathcal{F}_L \neq 0$.

Barlow's conjecture: $\mathcal{F}_{L+} \equiv \mathcal{F}_L \vee \sigma(A)$.

Proof of Lemma 2: It suffices to prove the 1st formula for elementary z 's: $z_u = 1_{\Gamma_t^1} 1_{]t, \infty]} (\frac{z}{u})$.

then:

$$\begin{aligned} E[z_L] &= E[1_{\Gamma_t^1} z_t^L] = -E[1_{\Gamma_t^1} (z_\infty^L - z_t^L)] \\ &= E[1_{\Gamma_t^1} (A_\infty^L - A_t^L)] \\ &= E\left[\int_0^\infty dA_u^L z_u\right]. \end{aligned}$$

For the 2nd formula, we use the first result, and integration by parts.

Proof of the theorem: We want to show:

$$E\left[\int_0^\infty H_u dX_u\right] = E\left[\int_0^\infty H_u \left\{ \underbrace{\frac{d\langle X, z \rangle_u^L}{z_u^L}}_{1_{(u < L)}} + \underbrace{\frac{d\langle X, 1-z \rangle_u^L}{(1-z_u^L)}}_{1_{(L < u)}} \right\}\right]$$

when (H_u) is a simple (\mathcal{F}_u^L) pred. process, i.e.: H' & H'' are simple.

$$\begin{aligned} &\text{We start from the RHS} \\ &= E\left[\int_0^L H'_u \frac{d\langle X, z \rangle_u^L}{z_u^L}\right] + E\left[\int_L^\infty H''_u \left(\frac{d\langle X, 1-z \rangle_u^L}{1-z_u^L}\right)\right] \\ &= E\left[\int_0^\infty H'_u d\langle X, z \rangle_u^L\right] + E\left[\int_0^\infty H''_u d\langle X, 1-z \rangle_u^L\right] \\ &= E\left[\int_0^\infty H'_u d\langle X, M \rangle_u^L\right] - E\left[\int_0^\infty H''_u d\langle X, M \rangle_u^L\right] \end{aligned}$$

$$= E \left[\int_0^L H'_u dX_u - \int_0^L H''_u dX_u \right]$$

$$= E \left[\int_0^L H'_u dX_u + \int_L^\infty H''_u dX_u \right] = E \left[\int_0^\infty H_u dX_u \right].$$

We then consider our previous example with $L = q_1$, but now we ~~also~~ have 2 enlargements formulas:

1st enl. formula [in $(\mathcal{F}_t^{\text{prog}, q_1})$]:

/ Recall: $Z_u^{q_1} \equiv P(q_1 > u | \mathcal{F}_u)$
 $\equiv \frac{\Phi}{\Phi'} \left(\frac{|B_u|}{\sqrt{1-u}} \right)$

$$B_t = \beta_t^{(\text{prog})} + \int_0^{t \wedge q_1} \left(-\frac{\Phi'}{\Phi} \right) \left(\frac{|B_u|}{\sqrt{1-u}} \right) \frac{\text{sgn}(B_u) du}{\sqrt{1-u}}$$

$$\mathcal{F}_t^{\text{prog}, q_1} \subset \mathcal{F}_t^{\text{in}, q_1} \quad / \quad + \int_{q_1}^t \left(\frac{\Phi'}{\Phi} \right) \left(\frac{|B_u|}{\sqrt{1-u}} \right) \frac{\text{sgn}(B_u)}{\sqrt{1-u}}$$

2nd enl. formula: [in $(\mathcal{F}_t^{\text{in}, q_1})$]:

$$B_t = \beta_t^{(\text{in})} - \int_0^{t \wedge q_1} \frac{du B_u}{(q_1 - u)} + \int_{q_1}^t \left(\frac{\Phi'}{\Phi} \right) \left(\frac{|B_u|}{\sqrt{1-u}} \right) \frac{\text{sgn}(B_u)}{\sqrt{1-u}}$$

Thus, we obtain:

$$\beta_t^{(\text{in})} = \beta_t^{(\text{prog})} + \int_0^{t \wedge q_1} \frac{du B_u}{(q_1 - u)} - \int_0^{t \wedge q_1} \left(\frac{\Phi'}{\Phi} \right) \left(\frac{|B_u|}{\sqrt{1-u}} \right) \frac{\text{sgn}(B_u)}{\sqrt{1-u}}$$

Pour que ces calculs soient compatibles, il faut donc savoir montrer:

$$E \left[\frac{1(u < q_1)}{(q_1 - u)} \mid \mathcal{F}_u^{\text{prog}, q_1} \right] |B_u| = \left(\frac{\Phi'}{\Phi} \right) \left(\frac{|B_u|}{\sqrt{1-u}} \right) \frac{1}{\sqrt{1-u}}$$

ceci découle d'une formule explicite (Lemme 3, de la 1^{re} exposé).

Details.]

un exemple, avec dernier temps de passage:

$(t, t \geq 0)$ (F_t) local mart; ≥ 0 ; $Y_t \xrightarrow[t \rightarrow \infty]{} 0$;

$< Y_0$; $\gamma_{xy}^d = \sup\{u \geq 0: Y_u = y\}$.

we have the following

Lemma:

i) $\sup_{t \geq 0} Y_t \stackrel{\text{law}}{=} \frac{Y_0}{U}$

ii) $Z_t^y \equiv P(\gamma_y^d > u | F_u) = \left(\frac{Y_u}{y}\right) \wedge 1.$

$\equiv Z_0^y + \frac{1}{y} \int_0^u 1(Y_v \leq y) dY_v - \frac{1}{2y} L_t^y.$

the enlarged formula becomes:

$$X_t = \tilde{X}_t + \int_0^{t \wedge \gamma_y^d} \frac{d\langle X, Y \rangle_v}{Y_v} 1(Y_v \leq y) - \int_0^t \frac{d\langle X, Y \rangle_v}{(y - Y_v)}$$

In particular, $\underbrace{Y_{y+v} - y}_{\leq 0} = \left(\tilde{Y}_{y+v} - \tilde{Y}_{y+y} \right) - \int_0^{t+y} \frac{d < Y_{v+y} >}{(y - Y_{v+y})}$

donc: 3 dim ?? BES /
 semble bizarre ... / (Vérifier)

En H cas, lorsque l'on applique ceci à $Y_t = 1/R_t$, avec $R = \text{BES}(3)$
 il vient: $\{ R_{y_1+t} - 1 ; t \geq 0 \} : \text{BES}(3)$.

A third note for Prof. Fujita.

At the end of the Pitman-Yor paper: "on the lengths of excursions...." (Sem. Prob. XXXI), we have a discussion of generalized arc sine laws, which may be applied again to yield some description of the joint law of (X_t, A_t^+) ,

where (X_t) is a skew Bessel process, with dimension $2-2\alpha$, and skewness parameter μ ; we denote the law of this process by $P_{\alpha, \mu}$. Then, we can write the following:

$$E_{\alpha, \mu} \left[f(X_t) \frac{1}{(1 + \lambda A_t^+)^{\alpha}} \mathbb{1}_{(X_t > 0)} \right] \stackrel{\text{def}}{=} \pi_{\alpha, \mu}(f, \lambda).$$

$$= E_{\alpha, \mu} \left[f(\sqrt{1-q} m_1) \frac{1}{(1 + \lambda [q A_{br}^+ + (1-q)])^{\alpha}} \right] \mu$$

where A_{br}^+ denotes the time spent ≥ 0 by the bridge.

Then, we also know the following:

- i) q is beta($\alpha, 1-\alpha$) distributed, this result goes back to Dynkin (1961)
- ii) the law of m_1 does not depend on α (this is a remarkable feature), i.e.: $P_{\alpha}(m_1 \in dp) = p e^{-p^2/2} dp.$
- iii) for every $\mu \geq 0$, $E_{\alpha, \mu} \left[\frac{1}{(1 + \mu A_{br}^+)^{\alpha}} \right] = \frac{1}{\mu(1+\mu)^{\alpha} + q}.$

If we write: $\pi_{\alpha, p}(f, \lambda) = p E_{\alpha, p} \left[f(\sqrt{1-q} m_1) \frac{1}{D^\alpha} \right]$,

where:

$$D = 1 + \lambda(1-q) + \lambda q A_{br}^+ \equiv (1 + \lambda(1-q)) (1 + \mu(q) A_{br}^+),$$

where: $\mu(q) \equiv \frac{\lambda q}{1 + \lambda(1-q)}$, we now deduce from (iii), that:

$$\pi_{\alpha, p}(f, \lambda) = E_{\alpha, p} \left[f(\sqrt{1-q} m_1) \frac{p}{(1 + \lambda(1-q))^\alpha [p(1 + \mu(q))^\alpha + q]} \right]$$

Since:

$$1 + \mu(q) \equiv \frac{1 + \lambda}{1 + \lambda(1-q)}, \text{ we finally obtain:}$$

$$\pi_{\alpha, p}(f, \lambda) = E_{\alpha, p} \left[f(\sqrt{1-q} m_1) \frac{p}{p(1 + \lambda)^\alpha + q(1 + \lambda(1-q))^\alpha} \right]$$

Now, it seems plausible that, for $a > 0$, there exists a n.v. $X_{p, a}$ such that:

$$E \left[\frac{1}{(1 + \lambda X_{p, a})^\alpha} \right] = \frac{1}{p(1 + \lambda)^\alpha + q(1 + \lambda a)^\alpha}$$

Assuming this is the case, we have obtained:

$$\pi_{\alpha, p}(f, \lambda) = E_{\alpha, p} \left[f(\sqrt{1-q} m_1) \frac{p}{(1 + \lambda X_{p, 1-q})^\alpha} \right],$$

and finally, we find that, conditionally on $(X_1 > 0)$,

$$(X_1, A_1^+) \stackrel{\text{(law)}}{=} (\sqrt{1-q} m_1, X_{p, 1-q}).$$