

(4)

On ranked lengths of excursions, III: from Brownian motion to the Brownian bridge.

July, 27<sup>th</sup>, 1994.

The main relation, which will allow us to pass to the Brownian bridge, is the following: in the sequel, let  $c = \sqrt{\frac{\pi}{2}}$ , and define:  $\beta_u^{(1)} \equiv \frac{1}{\sqrt{g_{H(1)}}} B_{u g_{H(1)}}$

Then, we have:

$$(1) \int c E \left[ F \left( \beta_u^{(1)}, u \leq 1 \right) f \left( g_{H(1)} \right) \right] = E \left[ \frac{1}{\sqrt{z_1}} F \left( \frac{B_{u z_1}}{\sqrt{z_1}}, u \leq 1 \right) \hat{f} \left( \frac{z_1}{V_1(z_1)} \right) \right]$$

$$\text{where: } \hat{f}(x) = \int_0^x \frac{du}{2\sqrt{u}} f(u)$$

and we can go from the pseudo-bridge  $B_u^\# \equiv \frac{1}{\sqrt{z_1}} B_{u z_1}$  to the bridge  $(b(u), u \leq 1)$  thanks to BLY, which gives:

$$(2) \int c^2 E \left[ F \left( \beta_u^{(1)}, u \leq 1 \right) f \left( g_{H(1)} \right) \right] = E \left[ F(b(u), u \leq 1) \hat{f} \left( \frac{1}{V_1(b)} \right) \right]$$

We now take  $f \equiv 1$ , and we get:

$$(1') \int_{f=1} c E \left[ F \left( \beta_u^{(1)}, u \leq 1 \right) \right] = E \left[ F \left( B_u^\#, u \leq 1 \right) \frac{1}{\sqrt{V_1(z_1)}} \right]$$

and

$$(2') \int_{f=1} c^2 E \left[ F \left( \beta_u^{(1)}, u \leq 1 \right) \right] = E \left[ F(b(u), u \leq 1) \frac{1}{\sqrt{V_1(b)}} \right].$$

Now, we apply  $(2')_{f=1}$  to  $F(X) = f\left(\frac{V_2(X)}{V_1(X)}, \frac{V_3(X)}{V_1(X)}, \dots\right)$  and we use two facts:

- a)  $V_{m+1}(H^{(1)}) = V_{m+1}(H^{(1)})$ ;
- b)  $H^{(1)}$  is admissible.

Thus, we deduce from  $(2')_{f=1}$  that:

$$(3) \quad \frac{\pi}{2} E\left[f\left(\frac{V_3(1)}{V_2(1)}, \frac{V_4(1)}{V_2(1)}, \dots\right)\right] = E\left[\frac{1}{\sqrt{V_1(b)}} f\left(\frac{V_2(b)}{V_1(b)}, \frac{V_3(b)}{V_1(b)}, \dots\right)\right]$$

Now, we remark that, if we denote:  $R_1^b = \frac{V_2(b)}{V_1(b)}$ ,  $R_m^b = \frac{V_{m+1}(b)}{V_m(b)}$ , then:

$$\frac{1}{\sqrt{V_1(b)}} = 1 + R_1^b + R_1^b R_2^b + R_1^b R_2^b R_3^b + \dots,$$

so that we can write (3) in the form:

$$(3') \quad \frac{\pi}{2} E\left[f(R_2, R_3, \dots)\right] = E\left[\left(1 + R_1^b + R_1^b R_2^b + \dots\right)^{1/2} f(R_1^b, R_2^b, \dots)\right]$$

or equivalently:

$$(3'') \quad E\left[f(R_1^b, R_2^b, \dots)\right] = \frac{\pi}{2} E\left[\left(1 + R_2 + R_2 R_3 + \dots\right)^{1/2} f(R_2, R_3, \dots)\right]$$

We now complete this discussion with remarks about the sequence:

$$\left(V_1(\bar{z}_1), V_2(\bar{z}_1), \dots\right)$$

Starting now from  $(1')_{f=1}$ , we obtain:

$$(4) \quad c E \left[ f \left( \frac{V_3(1)}{V_2(1)}, \frac{V_4(1)}{V_3(1)}, \frac{V_5(1)}{V_4(1)}, \dots \right) \right] = E \left[ \frac{1}{\sqrt{V_1(\bar{z}_1)}} f \left( \frac{V_2(\bar{z}_1)}{V_1(\bar{z}_1)}, \frac{V_3(\bar{z}_1)}{V_2(\bar{z}_1)}, \frac{V_4(\bar{z}_1)}{V_3(\bar{z}_1)}, \dots \right) \right]$$

(Note that, in order to deduce (4) from  $(1')_{f=1}$ , I have used again the facts a) and b) above).

Now, I introduce the notation:  $R'_1 = \frac{V_2(\bar{z}_1)}{V_1(\bar{z}_1)}$ ,  $R'_2 = \frac{V_3(\bar{z}_1)}{V_2(\bar{z}_1)}$ ,  $R'_3 = \frac{V_4(\bar{z}_1)}{V_3(\bar{z}_1)}$ , ...

We know, since  $\bar{z}_1$  is admissible, that:

$$(5) \quad (R_1, R_2, \dots, R_n, \dots) \stackrel{\text{law}}{=} (R'_1, R'_2, \dots, R'_n, \dots)$$

This being recalled, we see that formula (4) offers (... gives / yields ...) an expression for the Radon-Nikodym density of the distribution of:  $(R_{n+1}, n \geq 1)$  with respect to that of  $(R_n, n \geq 1)$ .

Since the  $R'_k$ 's are independent, it is easy to find an expression for this density; we shall do this now, and then compare with (4).

Proposition

1. For every  $n \geq 1$ , one has:

$$E \left[ f(R_2, R_3, \dots, R_{n+1}) \right] = E \left[ (n+1) (R_1 R_2 \dots R_n)^{1/2} f(R_1, R_2, \dots, R_n) \right]$$

2.  $(\Delta_n \stackrel{\text{def}}{=} (n+1) (R_1 R_2 \dots R_n)^{1/2}, n \geq 1)$  is an  $L^2$  convergent martingale, and we have:

for every measurable functional  $f$  defined on  $[0, 1]^{\mathbb{N}}$ ,  $f \geq 0$ ,

$$(6) \quad E \left[ f((R_{n+1}, n \geq 1)) \right] = E \left[ \Delta_\infty f((R_n, n \geq 1)) \right],$$

where:  $\Delta_\infty = \lim_{n \rightarrow \infty} n (R_1 R_2 \dots R_n)^{1/2}$ .

Proof: Recall that  $R_k \stackrel{\text{(law)}}{=} Z_{\frac{k}{2}, 1}$ ; hence, we have:

$$E[f(R_{k+1})] = \left(\frac{k+1}{2}\right) \int_0^1 dx x^{\frac{k+1}{2}-1} f(x) = \left(\frac{k+1}{k}\right) \left(\frac{k}{2}\right) \int_0^1 dx (x^{1/2}) x^{\frac{k}{2}-1} f(x)$$

$$= E\left[\left(\frac{k+1}{k}\right) (R_k)^{1/2} f(R_k)\right],$$

and the assertions in the proposition follow from the independence of the  $R_k$ 's  $\square$

We now compare (4) and (6): from the expression of  $\Delta_\infty$  given in the Proposition, we deduce that:

$$\frac{1}{c \sqrt{V_1(z_1)}} = \lim_{n \rightarrow \infty} n \left( \frac{V_2(z_1)}{V_1(z_1)} \cdot \frac{V_3(z_1)}{V_2(z_1)} \cdots \frac{V_{n+1}(z_1)}{V_n(z_1)} \right)$$

which simplifies, to give:  $\lim_{n \rightarrow \infty} c n (V_n(z_1))^{1/2} = 1$ ,

or equivalently: (7)  $\lim_{n \rightarrow \infty} \left( \frac{\pi}{2} n^2 V_n(z_1) \right) = 1$

This agrees with our previous limit theorem which I presented partly in my Fax of June 28<sup>th</sup> [and which you amplified, correcting my completely wrong conjecture!!] i.e.:

(8)  $\lim_{n \rightarrow \infty} \left( \frac{\pi}{2} n^2 V_n(t) \right) = L_t^2$

For clarity, I have now written 4 "Berkeley notes":

- (1) On ranked lengths of excursions, I
- (2) \_\_\_\_\_, II
- (3) Random scaling and admissibility
- (4) On ranked lengths of excursions, III:  
from Brownian motion to the Brownian bridge

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